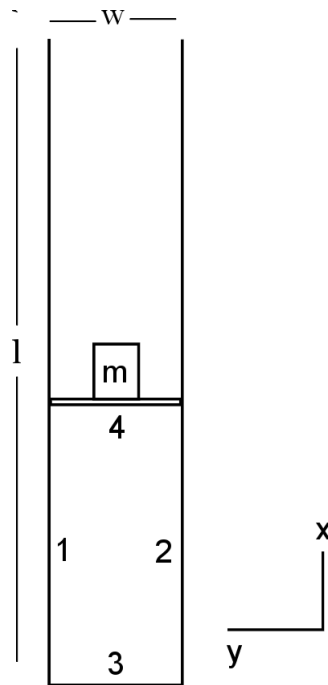


# Rail Gun

January 30, 2015

## Description of problem:

A projectile of mass  $m$  will be launched from a rail gun of length  $l$ , width  $w$ , resistance  $R$ . The power supply will be a capacitor bank of capacitance  $C$  at initial voltage of  $V_{\text{cap}}$ .



## Application of the Biot-Savart Law:

The magnetic field generated at any given point  $(x, y)$  (the mass's location) is given by the sum of the three other rail's (1, 2, and 3) contributions.

$$\vec{B} = \vec{B}_1 + \vec{B}_2 + \vec{B}_3$$

where

$$\vec{B}_1 = \frac{\mu_0 I}{4\pi y} \left( \frac{x}{\sqrt{x^2 + y^2}} + \frac{l-x}{\sqrt{(l-x)^2 + y^2}} \right) \hat{z}$$

$$\vec{B}_2 = \frac{\mu_0 I}{4\pi (w-y)} \left( \frac{l-x}{\sqrt{(l-x)^2 + (w-y)^2}} + \frac{x}{\sqrt{x^2 + (w-y)^2}} \right) \hat{z}$$

$$\vec{B}_3 = \frac{\mu_0 I}{4\pi x} \left( \frac{w-y}{\sqrt{x^2 + (w-y)^2}} + \frac{y}{\sqrt{x^2 + y^2}} \right) \hat{z}$$

For the electromotive force we include the field from the sliding current element as well

$$\vec{B}_4 = \frac{\mu_0 I}{4\pi (l-x)} \left( \frac{y}{\sqrt{y^2 + (l-x)^2}} + \frac{w-y}{\sqrt{(w-y)^2 + (l-x)^2}} \right) \hat{z}$$

so this defines the total as

$$\vec{B}_{\text{tot}} = \vec{B}_1 + \vec{B}_2 + \vec{B}_3 + \vec{B}_4$$

## Results of Induced EMF in loop:

Faraday tells us,

$$\nabla \times \vec{E} = -\frac{d\vec{B}_{\text{tot}}}{dt}$$

Using Stokes theorem

$$\int (\nabla \times \vec{E}) \cdot d\vec{A} = \oint \vec{E} \cdot d\vec{l}$$

we may write

$$\int \left( -\frac{d\vec{B}_{\text{tot}}}{dt} \right) \cdot d\vec{A} = \varepsilon \rightarrow \boxed{\varepsilon(t) = - \int \frac{d\vec{B}_{\text{tot}}}{dt} dx dy}$$

This induced EMF opposes the power supply which can be written in Ohm's law

$$V(t) = I(t) R = V_{\text{cap}}(t) - \varepsilon(t)$$

Therefore, the current that can be used in Biot-Savart application is

$$I(t) = \frac{V_{\text{cap}}(t) - \varepsilon(t)}{R}$$

To add details, the power supply's voltage is given as

$$C = \frac{q(t)}{V_{\text{cap}}(t)} \rightarrow V_{\text{cap}}(t) = \frac{1}{C} \left[ q(0) - \int I(t) dt \right]$$

because discharging capacitors obey

$$q(t) = q(0) - \int I(t) dt$$

making the current

$$I(t) = \frac{1}{C} \frac{[q(0) - \int I(t) dt] - \varepsilon(t)}{R}$$

The time-derivative of this is

$$\frac{dI}{dt} = \frac{d}{dt} \left( \frac{1}{RC} \left[ q(0) - \int I(t) dt \right] - \frac{\varepsilon(t)}{R} \right) = \frac{1}{RC} \frac{d}{dt} q(0) - \frac{1}{RC} \frac{d}{dt} \int I(t) dt - \frac{d\varepsilon(t)}{dt} \frac{1}{R}$$

$$\boxed{\frac{dI}{dt} = -\frac{I(t)}{RC} - \frac{1}{R} \frac{d\varepsilon}{dt}}$$

### Force on the mass:

The Lorentz force is given by

$$\vec{F} = I(t) \int d\vec{l} \times \vec{B}_{\text{tot}} = I(t) \int_0^w |\vec{B}_{\text{tot}}| dy$$

$$\boxed{m \frac{d^2}{dt^2} x(t) = \frac{\mu_0 x(t) I^2(t)}{\pi} \left( \frac{w^2 + x^2(t) - x(t) \sqrt{w^2 + x^2(t)}}{\sqrt{w^2 + x^2(t)}} \right)}$$

and the equation for the current:

$$\boxed{\frac{dI}{dt} = -\frac{I(t)}{RC} + \frac{1}{R} \frac{d}{dt} \left( \int \frac{d\vec{B}_{\text{tot}}}{dt} dx dy \right)}$$

Now to solve these equations!